

AD A118894

AR-002-328



DEPARTMENT OF DEFENCE DEFENCE SCIENCE AND TECHNOLOGY ORGANISATION AERONAUTICAL RESEARCH LABORATORIES

MELBOURNE, VICTORIA

AERODYNAMICS NOTE 405

AN INTERFEROMETRIC INVESTIGATION OF THE NEAR DESIGN POINT FLOW OVER SUPERCRITICAL AEROFOIL BGK-1

by

N. POLLOCK

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AN INTERFEROMETRIC INVESTIGATION OF THE NEAR DESIGN POINT FLOW OVER SUPERCRITICAL AEROFOIL BGK-1

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SUMMARY

A recently developed laser interferometer was used to study the upper surface flow over supercritical aerofoil BGK-1 near its design condition. The tests were carried out at a chord Reynolds number of 1.65×10^6 with transition free and artificially fixed at various chordwise locations. The interferograms were analysed to produce "instantaneous" surface pressure distributions for comparison with conventional time averaged distributions obtained from surface tappings.

The results show that the upper surface flow is unsteady with both fixed and free transition. They also indicate that artificially fixing transition in low Reynolds number tests does not produce an accurate simulation of the high Reynolds number flow for conditions near the design point.



CONTENTS

NOTATION	Page No.
1. INTRODUCTION	1
2. EXPERIMENTAL DETAILS	1
2.1 Model	1
2.2 Tunnel	1
2.3 Instrumentation	2
2.4 Test program	2
3. RESULTS AND DISCUSSION	3
3.1 Transition Free, Design Conditions	3
3.2 Transition Free, Off-design Conditions	4
3.3 Transition Fixed	4
4. CONCLUSIONS	5
REFERENCES	
TABLES	
APPENDIX	
FIGURES	
DISTRIBUTION	

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NOTATION

- c Model chord length
- H Total pressure of tunnel flow
- M Mach number
- P Static pressure acting on model
- x Streamwise ordinate measured from leading edge
- y Thickness ordinate measured from chord line
- Model angle of incidence

1. INTRODUCTION

During the last 15 years many supercritical aerofoil sections have been designed and a number of these sections are now being used on production aircraft. Perhaps the most widely tested 1-7 of these sections is the profile designated BGK-1 designed by the numerical complex hodograph method of Bauer, Garabedian and Korn⁸. Although not intended to be a practical wing Section BGK-1 was designed to closely resemble some of the early semi-empirical NASA supercritical sections⁹. This similarity to practical empirical designs and the existence of an accurately known inviscid design pressure distribution explains the popularity of BGK-1 as a standard test aerofoil for wind tunnels and computational methods.

Some of the wind tunnel tests carried out on BGK-1 at a Reynolds number of 1.6×10^6 and without artificial transition fixing show a pressure distribution at the design point in strikingly good agreement with theoretical predictions⁸. Higher Reynolds number tests^{1,2} and tests with artificial transition fixing^{4,6} showed significantly inferior agreement with theory. The available surface pressure distributions and schlieren and shadowgraph optical flow visualisation photographs do not contain sufficient information to explain these differences.

During 1979 a new laser interferometer system¹⁰ was installed in the ARL transonic wind tunnel. It was decided to use the new flow visualisation capability provided by this system to further elucidate the nature of the near design point flow around supercritical aerofoil BGK-1. It was considered that the availability of interferograms of the flow would significantly increase the value of BGK-1 as a standard test aerofoil for both experimental and theoretical methods.

The tests described in this note were carried out in the transonic wind tunnel during September 1980.

2. EXPERIMENTAL DETAILS

2.1 Model

The BGK-1 aerofoil (Fig. 1) model used for these tests has a chord of $203 \cdot 2$ mm and a span of 532 mm. The design coordinates used for manufacturing were obtained from Reference 1. In the range $0.00 \le x/c \le 0.95$ the coordinates used for manufacturing were identical to the design coordinates (Table 1-1 and 1-2). For x/c > 0.95 the coordinates were modified to produce a trailing edge thickness of 0.0005c (Table 1-3). The measured profile errors for the completed model are plotted in Figure 2. On the more critical upper surface the accuracy achieved falls within the tolerances suggested in Reference 10 for this type of model (local waviness < 0.0002c and overall errors < 0.0005c).

The model was made from a number of pieces of steel joined rigidly prior to the final machining of the profile. This enabled the surface pressure hole connections to be made without spoiling the windswept surfaces. The model had 37 pressure holes 0.4 mm diameter in the upper surface (Table 2). The holes were distributed over a span wise region 0 19c wide centred at midspan. The model was also equipped with pressure holes in the lower surface but they were not used in these tests.

For the tests conducted with fixed transition spanwise roughness bands comprising carborundum particles attached with lacquer were used. The bands were 1.5%c wide with a roughness height of approximately 0.15 mm. The lacquer layer was 0.03 mm thick and the roughness coverage was 10% to 20%.

2.2 Tunnel

The test section used was 533 mm wide and 813 mm high (Fig. 3). Solid sidewalls and longitudinally slotted top and bottom walls (open area ratio of 16.5% at the model location) were fitted. Mach number was derived from measurements of the pressure in the plenum chamber

surrounding the test section and the contraction entry, assuming these to be the static and total pressures of the test section flow respectively.

The model was supported by means of integral end tongues fitting into slots in sidewall mounted optical glass windows. The pressure lines were brought through the tongues to the outside of the test section.

There were gaps of 0.5 mm between the ends of the aerofoil section and the glass windows. Brief tests with these gaps sealed indicated that their presence did not affect the results.

2.3 Instrumentation

The model surface pressures were measured using two "Scanivalve" pressure multiplexers fitted with ± 33.5 kPa differential pressure transducers referenced to tunnel static pressure. Stagger scanning was used to ensure the maximum settling time between readings. Scanning speed was 14 ports per second, with free stream conditions being recorded after each group of 12 ports. The model pressure tappings and connecting tube between model and scanivalve form a pneumatic low pass filter with a time constant of the order of 1 second, therefore pressure fluctuations at more than a few Hertz will not appear in the resulting pressure distributions.

The laser interferometer used is described in Reference 11. The field of view available with this instrument was slightly greater than half the area of the 406 mm diameter tunnel windows. The interferometer was therefore arranged so that the model upper surface was visible and the flow over the less interesting (for this type of aerofoil) lower surface was not recorded. The interferograms obtained were analysed to produce surface pressure distributions using the method described in the Appendix. The exposure time used in the interferometric investigation was $5 \,\mu \text{sec}$. so pressure fluctuations at up to about 100 kHz should be reproduced. A typical interferogram is reproduced in Figure 4. Since the fringes are difficult to follow in some areas the interferometric data have been provided in the form of fringe tracings (Figs 5–8 and 10–25) rather than photographic reproductions.

The arrangement of the interferometer (Ref. 11) is such that the effect of vibration is to uniformly stretch or compress the fringe pattern in a direction normal to the aerofoil chord. Empty tunnel tests indicate that this effect is less than one fringe over the whole field. There is no mechanism for mechanical vibration to produce local variations in the fringe pattern. Previous tests with the interferometer have confirmed that accurate, reproducible interferograms are obtained when the flow is known to be steady.

When comparing surface pressure measurements with pressure distributions reduced from interferograms it should be remembered that the former are measured near midspan and the latter are a full span average (see Appendix). However at the low incidence angles of these tests the flow should be highly two-dimensional and the difference should not materially affect comparisons between the two measurements. There is a possibility that unsteadiness confined to the tunnel side wall boundary layers will be visible in the interferograms, but since the wall boundary layer thickness is small compared to the tun width this should not be a significant problem.

2.4 Test Program

Previous tests⁸ on this aerofoil in the same wind tunnel have shown that the experimental design point (i.e. the test conditions which best reproduce the theoretical upper surface design pressure distribution) occurs at M = 0.783 and $\alpha = 1.7^{\circ}$. The present test program was therefore centered on these values.

The difference between the experimental design conditions M=0.783, $\alpha=1.7^{\circ}$ and the theoretical design values M=0.75, $\alpha=0^{\circ}$ is due partly to the effects of viscosity and partly due to tunnel interference. When corrected for tunnel interference the experimental design conditions become M=0.760, $\alpha=0.50^{\circ}$.

The chord Reynolds number for this series of tests was $1.65 \pm 0.05 \times 10^6$. Several interferograms were produced at some test conditions to provide information on flow steadiness. (The particular cases tested are listed in the table below:)

М	α	Transition	Figure Numbers
0.783	1 · 7°	Free	4_9
0.775	1 · 7°	Free	10
0.780	1 · 7°	Free	11
0.788	1 · 7°	Free	12
0.793	1 · 7°	Free	13
0.783	1 · 2°	Free	14
0.783	1 · 5°	Free	15
0.783	1 · 9°	Free	16
0.783	2·1°	Free	17
0.783	1 · 7°	Fixed at 2%c	18–20
0.783	1 · 7°	Fixed at 5%c	21
0.783	1 · 7°	Fixed at 17%c	22 and 23
0.783	1 · 7°	Fixed at 37%c	24 and 25

3. RESULTS AND DISCUSSION

3.1 Transition Free, Design Conditions

In Figures 5 to 8 four interferograms of the experimental design point flow $(M=0.783, \alpha=1.7^{\circ})$ with transition free are presented. Pressure distributions reduced from the interferograms are also presented along with measured surface pressure distributions and the design pressure distribution. The four measured surface distributions are very nearly identical and in good agreement with the theoretical distribution. Previous work⁸ has shown that M and α increments of 0.002 and 0.1° respectively, markedly reduced the agreement with the theoretical distribution.

Unlike the measured surface distributions the pressure distributions reduced from the interferograms showed considerable variations particularly near the end of the supersonic region at 0.4 < x/c < 0.7. The interferograms (Fig. 4) and schlieren flow visualisation (Fig. 9) indicate that boundary layer transition takes place in this x/c range. The interferograms also show that away from the model surface the initial expansion near the leading edge (fringes sloping upwards to the right) is followed by a region of recompression (fringes sloping downwards to the right). This recompression terminates in the range 0.4 < x/c < 0.5 and is followed by another expansion and compression. In some cases (e.g. Fig. 8) further weaker expansions and compressions occur prior to the final smooth subsonic recompression approaching the trailing edge. This cyclic expansion and compression is also evident in the long exposure schlieren photograph in Figure 9. Some of the interferograms (e.g. Fig. 7) show lines formed by the confluence of fringes indicating shock waves. Similar shock waves are evident in the short exposure schlieren photograph in Figure 9. Two of the reduced interferogram pressure distributions (Figs 5 and 8) show a pressure plateau typical of a shock induced laminar boundary layer separation followed by a turbulent reattachment 14. The two others (Figs 6 and 7) do not show a clear plateau but the highly curved fringes in the interferograms near the model surface in the range 0.5 < x/c < 0.7 indicates that the boundary layer displacement surface is significantly distorted.

The conclusion to be drawn from the above observations is that an unsteady laminar/ transitional shock wave boundary layer interaction occurs in the range 0.5 < x/c < 0.7. The test Reynolds number appears to be such that a significant laminar separation bubble is present for some of the time only. The significance of the fact that the time averaged pressure distribution (as measured by the surface tappings) is smooth and in good agreement with the theoretical distribution is not clear.

It is suggested that this design point flow may be useful as a particularly difficult test case for numerical Navier-Stokes equation solutions. The strong link between the boundary layer transition process and the outer flow would place particular emphasis on the correct representation of the physics of the flow in the numerical formulation.

3.2 Transition Free, Off Design Conditions

At Mach numbers below the design value and at the design incidence (Figs 10 and 11) both the measured surface pressures and the reduced interferograms show that the initial supersonic recompression is somewhat steeper than the design case with a significant reexpansion preceding the final recompression to subsonic speed. At Mach numbers above the design value (Figs 12 and 13) the initial supersonic recompression tends to develop into a pressure plateau with a final steep recompression to subsonic speed.

At incidence angles below the design value and at the design Mach number (Figs 14 and 15) the initial recompression is steeper than the design case with a pressure plateau or very weak expansion preceding the final return to subsonic speed. At incidence angles above the design value (Figs 16 and 17) the initial recompression is weakened with a rapid final return to subsonic speed involving shock waves.

The differences between the measured surface pressures and those reduced from the interferograms indicates that, as for the design point tests, the flow is quite unsteady under close to design conditions. Comparing the off-design (Figs 10-17) and design point (Figs 5-8) results it is evident that the time averaged design distribution evolves smoothly from the off-design distributions as the values of M and α are varied. As noted previously very small increments of M and α produce significant changes in the time averaged pressure distributions. The interferograms and the "instantaneous" pressure distributions derived from them show considerable variation and it would be difficult to distinguish between design and off-design conditions on this basis alone (see for example Figs 7 and 10). This again raises the question of the significance of the time averaged pressure distribution in an unsteady flow. It is tempting to consider the time averaged flow as the underlying steady flow on which are superimposed random fluctuations caused by tunnel flow unsteadiness. The alternative possibility is that the flow is inherently unstable and that the time averaged flow is not in itself a legitimate or possible flow. Numerical calculations would be valuable in resolving this question.

The off-design results (Figs 10 to 17) show that boundary layer transition (indicated by the appearance of the broken line showing the visible edge of the turbulent boundary layer) and the termination of the initial smooth supersonic recompression are closely associated. On the evidence available it is not possible to state which one causes the other, if indeed they do have a cause and effect relationship. Numerical calculations may also be useful in answering this question.

3.3 Transition Fixed

Results of tests with transition fixed at 2%c are presented in Figures 18-20. This transition location is in the subsonic part of the leading edge expansion and does not in itself produce a major impact on the pressure distribution or the interferogram. The resulting pressure distributions and interferograms have some similarity with the transition free results at less than the design Mach number (Fig. 10). However previous tests⁴ have shown that the agreement between the transition fixed results and the theoretical design pressure distribution is not improved by varying the Mach number or incidence. The most probable mechanism for transition fixing to alter the design point pressure distribution is the alteration of the effective section displacement surface by the relatively thick turbulent boundary layer. This is consistent with the known sensitivity of this type of aerofoil to small changes in surface shape. Comparing Figures 18-20 it is evident that the flow has significant unsteadiness. This observation suggests that the unsteadiness noted in the transition free tests was not caused solely by the presence of an unsteady laminar boundary layer separation.

Moving the transition strip aft to 5% c (Fig. 21) does not appear to have any fundamental effect on the nature of the flow. Moving the transition strip further aft to 17% c where it is well into the supersonic region (Figs 22, 23) produces a large disturbance which extends well out into the flow. Despite this disturbance, which is clearly visible in the interferograms, the agreement between the surface pressure distributions and the design distribution is not significantly degraded over the forward transition cases discussed previously. This flow disturbance is believed to be due to the change in slope of the boundary layer displacement surface following transition rather than the obstruction produced by the transition strip itself. This view is supported by

the inviseid calculations of the effect of roughness bands presented in Reference 4. Moving the roughness band back to 37% c (Figs 24, 25) simply moves the flow disturbance aft. The results with transition at 17% c and 37% c (Figs 22-25) all show evidence of significant flow unsteadiness.

4. CONCLUSION

A recently developed laser interferometer has been used to study the upper surface flow over supercritical aerofoil BGK-1 near its design condition. The interferograms have been analysed to produce "instantaneous surface pressure distributions" which are compared with conventional time averaged pressure distributions obtained from surface tappings. These new data should considerably enhance the value of BGK-1 as a standard test aerofoil for experimental and theoretical methods.

The smooth "design" pressure distribution noted in previous transition free tests was found to be the time average of a very unsteady flow associated with unsteady boundary layer transition. It is not clear from the evidence available whether a laminar separation bubble is present some or all of the time at the design point. The significance of the excellent agreement between the theoretical design pressure distribution and the time average of a very unsteady flow appears to warrant further investigation.

Tests with boundary layer transition fixed with roughness bands indicated that transition locations in the supersonic region caused large distortions of the flow field for a significant distance from the model. Transition locations in the subsonic expansion near the nose did not obviously distort the flow, but the resulting thick turbulent boundary layer did destroy the agreement between theoretical and experimental design point pressure distributions. All the transition fixed tests showed evidence of flow unsteadiness which suggests that the transition free unsteadiness was not caused simply by unsteady transition location or laminar separation. It appears that, for this type of section, artificial transition fixing cannot be used to accurately simulate high Reynolds number conditions near the design point in low Reynolds number tunnel tests.

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TABLE 1-1
Profile of Aerofoil

Upper Surface—Design and Model $0.00 \le x/c \le 0.95$

x c	<u>y</u>	x -	$\frac{y}{c}$	x c	y
<u>с</u>	c	с	c	c	y c
0.00000	0.00219	0.07532	0.04352	0.27237	0.06271
0.00001	0.00298	0.07597	0.04366	0.28741	0.06336
0.00006	0.00403	0.07660	0.04379	0.30299	0.06396
0.00017	0.00528	0.07724	0.04392	0.31906	0.06450
0.00027	0.00603	0.07786	0.04405	0.33556	0.06496
0.00039	0.00673	0.07849	0.04418	0.35232	0.06535
0.00058	0.00758	0.07911	0.04430	0.36959	0.06566
0.00080	0.00832	0.08045	0.04457	0.38700	0.06588
0.00109	0.00913	0.07978	0.04444	0.40413	0.06602
0.00141	0.00987	0.08253	0.04496	0.42139	0.06608
0.00154	0.01015	0.08113	0.04470	0.43871	0.06604
0.00215	0.01124	0.08182	0.04483	0.45601	0.06592
0.00269	0.01203	0.08327	0.04510	0.47325	0.06572
0.00333	0.01286	0.08404	0.04524	0.49034	0.06542
0.00428	0.01392	0.08484	0.04539	0.50724	0.06504
0.00534	0.01498	0.08574	0.04555	0.52388	0.06457
0.00689	0.01637	0.08670	0.04572	0.53954	0.06405
0.00865	0.01778	0.08773	0.04589	0.55487	0.06345
0.01107	0.01953	0.08885	0.04608	0.56982	0.06478
0.01458	0.02178	0.09006	0.04628	0.58435	0.06203
0·01964	0.02464	0.09138	0.04650	0·59841	0.06122
0.02348	0.02658	0.09282	0.04674	0.61194	0.06035
0.02831	0.02881	0.09441	0.04699	0 62490	0.05943
0.03389	0.03116	0.09626	0.04728	0.63727	0.05847
0.03779	0.03267	0.09831	0.04759	0.64844	0.05754
0.04089	0.03380	0.10058	0.04793	0.65902	0.05659
0.05268	0.03766	0.10310	0.04829	0.66899	0.05564
0.05397	0.03805	0.10591	0.04869	0.67832	0.05469
0.05543	0.03848	0.10901	0.04912	0.68701	0.05377
0.05692	0.03891	0.11245	0.04959	0.69509	0.05288
0.05769	0.03913	0.11625	0.05009	0 · 70694	0.05149
0.05846	0.03935	0.12062	0.05064	0.72841	0.04877
0.05922	0.03956	0-12545	0.05124	0.73865	0.04738
0.05997	0.03977	0·13078	0.05187	0.76139	0.04410
0.06071	0.03997	0.13664	0.05254	0.77701	0.04168
0.06145	0.04017	0 · 14305	0.05324	0.79224	0.03920
0.06217	0.04037	0.15004	0.05398	0.80907	0.03633
0.06291	0.04056	0.15763	0.05474	0.82499	0.03349
0.06444	0.04096	0.16585	0.05552	0.83695	0.03128
0.06593	0.04133	0.17493	0.05634	0.85719	0.02742
0.06738	0.04169	0 · 18470	0.05717	0.87562	0.02379
0.06879	0.04204	0.19518	0.05801	0.89827	0.01923
0.07016	0.04236	0 · 20636	0.05884	0.91369	0.01610
0.07148	0.04267	0-21825	0.05967	0.92769	0.01328
0.07277	0.04296	0.23081	0.06048	0.93981	0.01087
0.07401	0.04324	0 · 24405	0.06126		
0·07467	0.04338	0.25792	0.06201	}	

TABLE 1-2
Profile of Aerofoil
Lower Surface—Design and Model $0.00 \le x/c \le 0.95$

$0.00 \leqslant x/c \leqslant 0.95$						
	!	! x	, ,			
X .	y –		. y			
С	c	<i>c</i>	C			
0.00001	0.00149	0.26400	-0.05141			
0.00004	0.00080	0.28205	-0.05150			
0.00018	-0.00068	0.29705	-0.05144			
0.00025	-0.00114	0.31015	-0.05128			
0.00035	-0.00169	0.32408	-0.05103			
0.00046	-0.00219	0.34487	-0.05046			
0.00068	-0.00301	0.36350	-0.04976			
0.00101	-0.00395	0.38765	-0.04860			
0.00148	-0.00500	0.40449	-0.04763			
0.00188	-0.00575	0.42631	-0.04616			
0.00232	-0.00649	0.45002	-0.04432			
0.00306	0.00760	0.47538	-0.04206			
0.00420	-0.00903	0.52184	-0.03720			
0.00475	-0.00964	0.54984	-0.03383			
0.00578	-0.01069	0.57861	-0.03008			
0.00704	-0.01181	0.60531	-0.02636			
0.00827	-0.01279	0.62348	-0.02373			
0.00997	-0.01401	0.64215	-0.02096			
0.01229	-0.01547	0.65754	-0.01865			
0.01535	-0.01718	0.67121	-0.01659			
0.01928	0.01911	0.69034	-0.01373			
0.02534	-0.02170	0.70889	-0.01101			
0.02938	-0.02324	0.72669	-0.00850			
0.03910	-0.02653	0.74668	0.00585			
0.05252	-0.03034	0.76850	-0.00322			
0.06247	-0.03278	0.78489	-0.00148			
0.07333	-0.03514	0.80115	0.00006			
0.08529	-0.03744	0.81919	0.00149			
0.09982	-0.03988	0.83855	0.00273			
0.11856	-0.04253	0.85882	0.00369			
0.14099	-0.04511	0.87943	0.00431			
0.16719	-0.04745	0.90372	0.00459			
0.20293	-0.04966	0.92651	0.00439			
0.22851	-0.05068	0.94365	0.00395			
0.24790	-0.05117					

TABLE 1-3

Profile of Aerofoil $0.95 \le x/c \le 1.00$ Upper Surface

<u>x</u>	y	y
c	c	c
	Design	Model
0.95200	0.00849	0.00851
0.97179	0.00480	0.00488
0.98910	0.00182	0.00200
0.99786	0.00048	0.00071
1.00000	0.00018	0.00043

Lower Surface

x	y	y
c	c	c
	Design	Model
0.95867	0.00333	0.00327
0.97362	0.00247	0.00238
0.98619	0.00154	0.00138
0.99658	0.00056	0.00034
1.00000	0.00018	-0.00001

TABLE 2
Pressure Hole Locations

Uppe	r Surface	
x y		
c	c	
0.0000	0.0047	
0.0035	0.0129	
0.0079	0.0179	
0.0196	0.0247	
0.0294	0.0293	
0.0495	0.0367	
0.0741	0.0433	
0.0991	0.0480	
*0.1241	0.0513	
0.1591	0.0549	
0.1994	0.0583	
0.2382	0.0609	
0.2794	0.0631	
0.3188	0.0645	
0.3594	0.0655	
0.3993	0.0660	
0.4194	0.0660	
0.4395	0.0660	
0.4585	0.0659	
0.4787	0.0656	
0.4983	0.0653	
0.5184	0.0648	
0.5396	0.0641	
	0.0632	
	0.0624	
0.5989	0.0612	
0·6180 0·6484	0.0599	
0·6484 0·6747	0.0576	
0.6989	0.0551	
, '	0·0525 0·0494	
0 = 400	0.0494	
0.7739	0.0422	
0.7000	0.0422	
0.7990	0.0381	
0.8489	0.0339	
0.8989	0.0191	
0 0/0/	0.0131	

^{*} Damaged orifice.

APPENDIX

Interferogram Analysis

The basic expression for analysing interferograms of two dimensional isentropic flows¹² is:

$$\frac{P}{H} = \left[\left(\frac{Pr}{H} \right)^{1/\gamma} + \frac{\lambda RT_0}{KLH} \cdot N \right]^{\gamma}$$

where

P =Static pressure at any fringe

H = Tunnel stagnation pressure

Pr = Reference pressure (i.e. pressure at fringe "0")

 T_0 = Tunnel stagnation temperature

L = Effective tunnel width (with allowance for tunnel sidewall boundary layers)

R = Gas constant

K = Gladstone Dale constant

 λ = Wavelength of light used

 γ = Ratio of specific heats

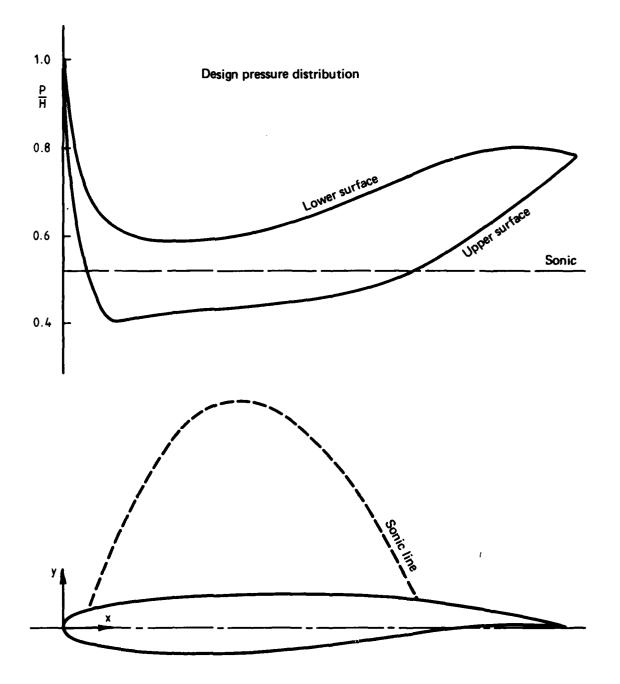
N =Fringe number where P is required.

For all the cases presented in this Note Fringe "0" was taken to be the first visible fringe to meet the model surface near the leading edge. The value of Pr for this fringe was obtained by interpolating the measured surface pressures.

The value of the parameter $\lambda R/KL$ was obtained by analysing the relatively steady monotonic compression which occurs over the forward 30% of chord. Comparing the number of fringes in this region with surface pressure distributions for five separate test cases produced a consistent value for $\lambda R/KL$. This procedure evaluates the otherwise unknown effective tunnel width.

The determination of the fringe number N is usually quite straightforward if the flowfield is worked through methodically starting with the reference fringe. However when shockwaves appear they tend to subdivide the flow into separate regions. Fortunately for all the cases presented here the shock waves were sufficiently short for an unbroken fringe to be found over the top of the shock effectively linking the two regions.

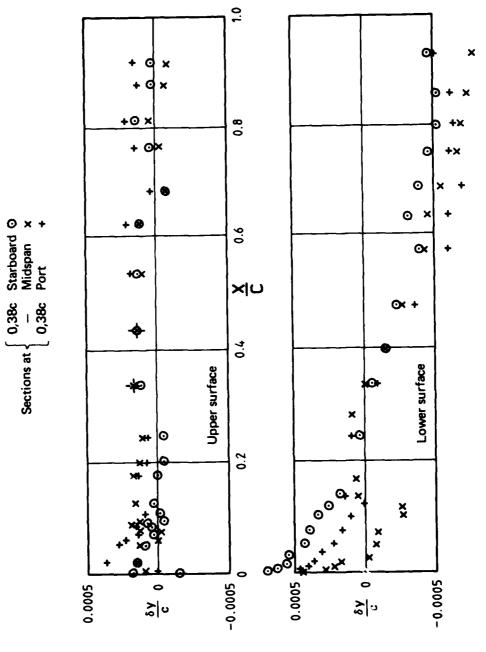
All the shock waves present in the tests reported here are considered to be sufficiently weak for the above isentropic analysis to give accurate results.



Thickness/Chord = 0,115

Design conditions M = 0.750 $\alpha = 0^{\circ}$ $C_{\perp} = 0.63$

FIG. 1 SUPERCRITICAL LIFTING AEROFOIL BGK1



δy = Measured ordinate — Manufacturing ordinate (positive for excess metal)

FIG. 2 PROFILE ERRORS

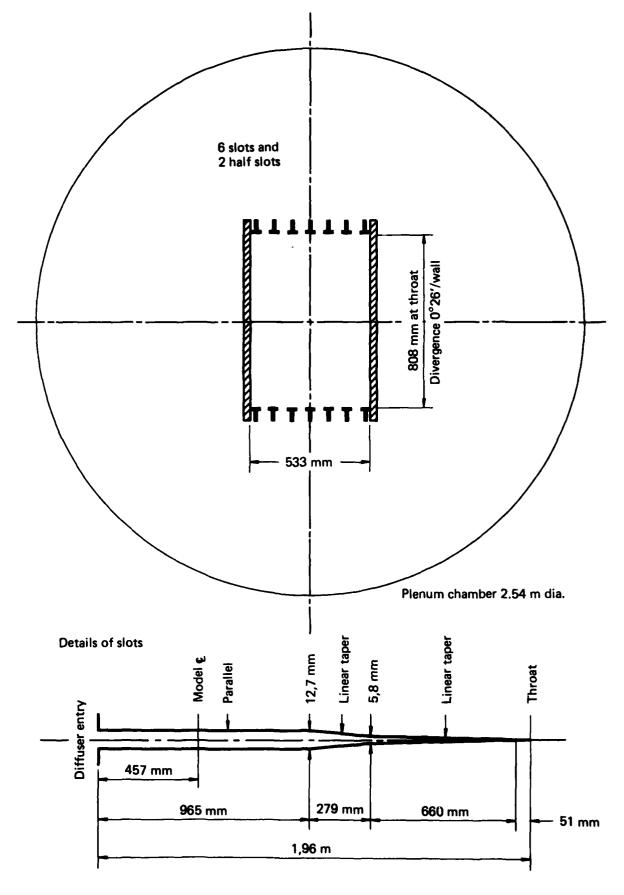


FIG. 3 DETAILS OF SLOTTED WORKING SECTION

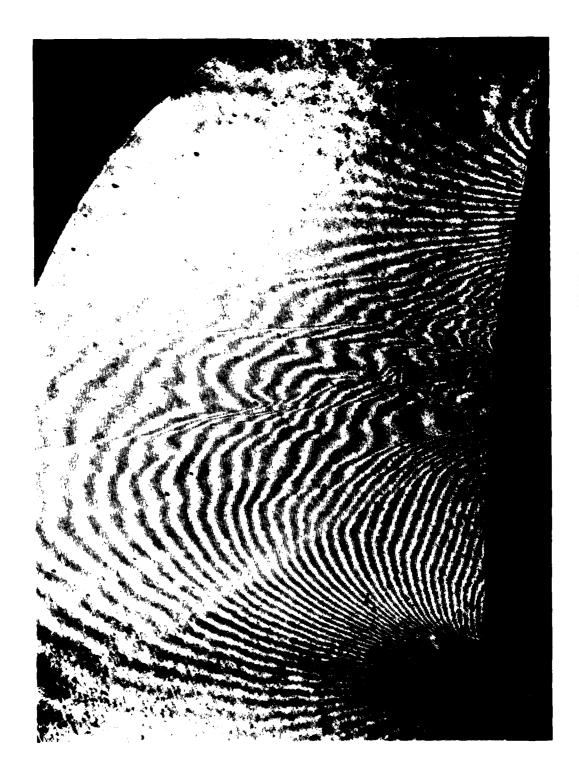
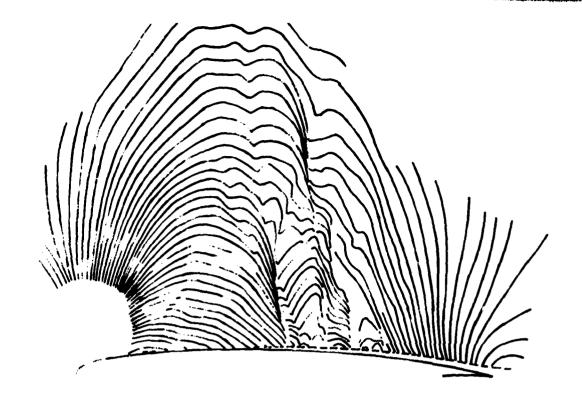
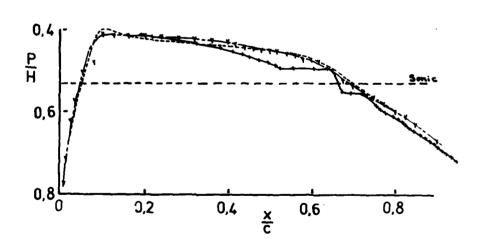


FIG. 4 PHOTOGRAPH OF INTERFEROGRAM $M=0.783, \alpha=1.7^{\circ}, TRANSITION FREE$





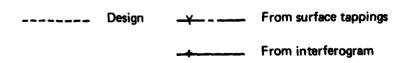
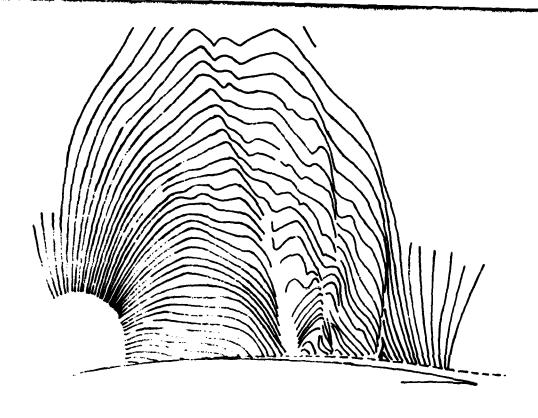
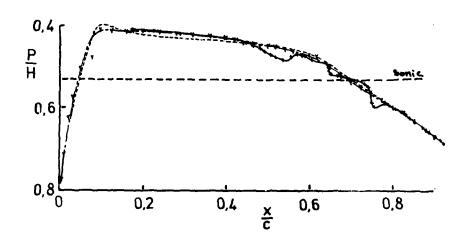


FIG. 5 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.7°, TRANSITION FREE





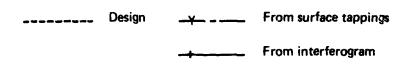
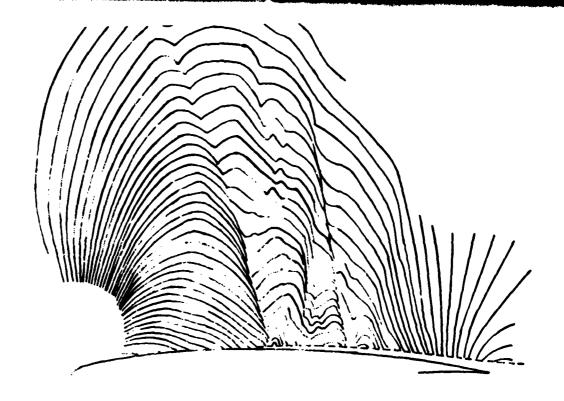
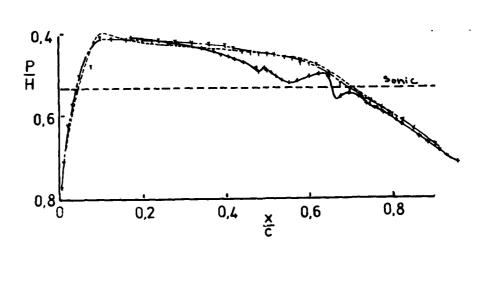


FIG. 6 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, $\alpha = 1.7^{\circ}$, TRANSITION FREE

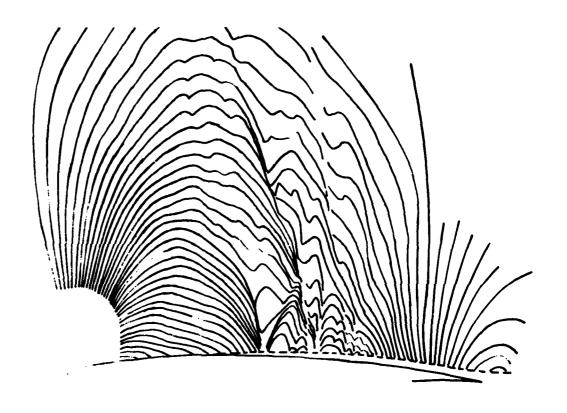


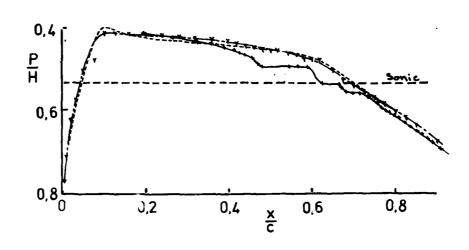


Design ——— From surface tappings

From interferogram

FIG. 7 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, $\alpha = 1.7^{\circ}$, TRANSITION FREE





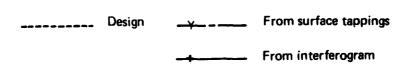
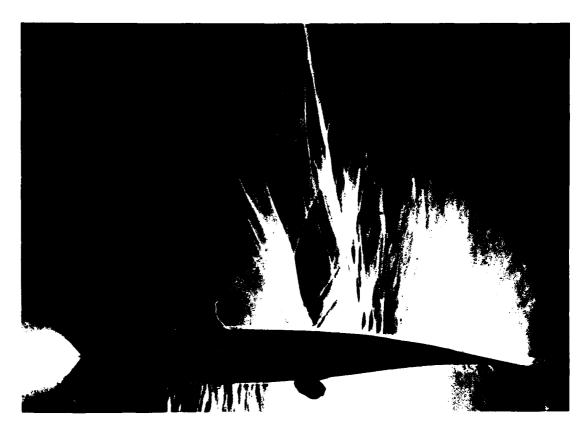
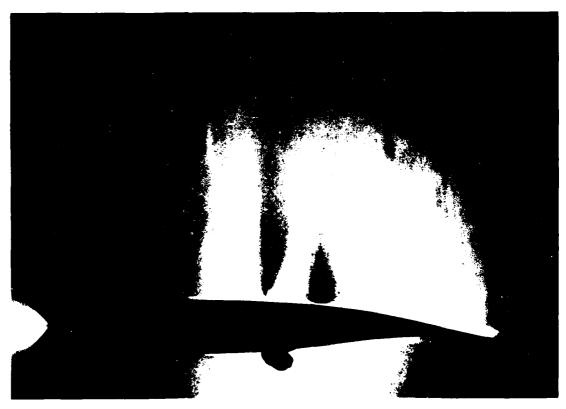


FIG. 8 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.7°, TRANSITION FREE

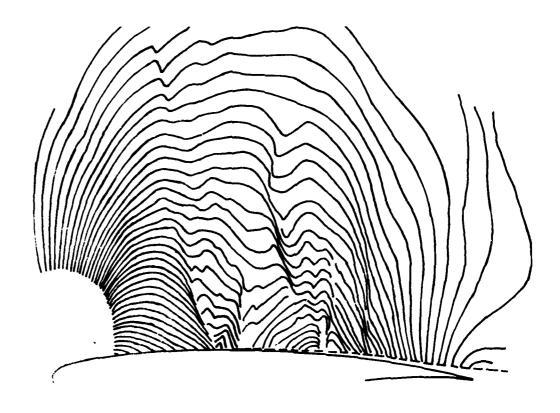


Exposure time = 1μ sec



Exposure time = 20 m sec

FIG. 9 SCHLIEREN PHOTOGRAPHS M = 0.783, α = 1.7°, TRANSITION FREE



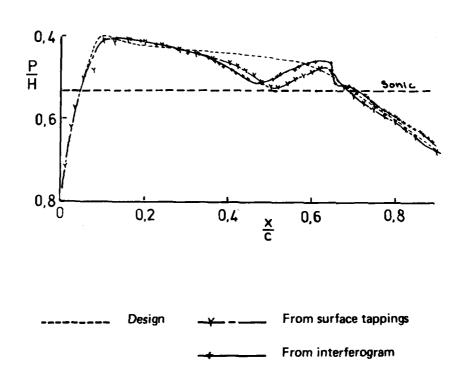
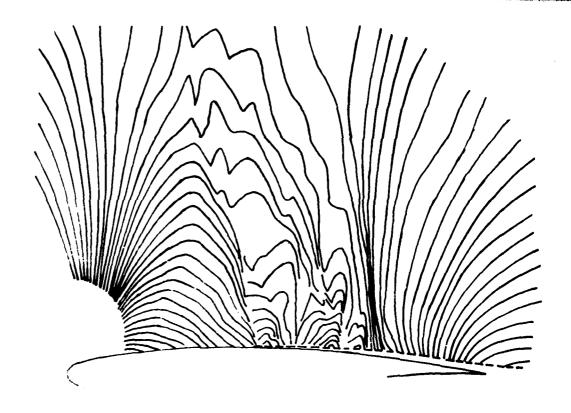
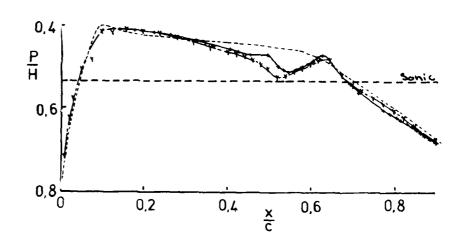


FIG. 10 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS $M=0.775, \alpha=1.7^{\circ}, TRANSITION FREE$





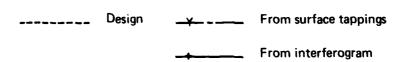
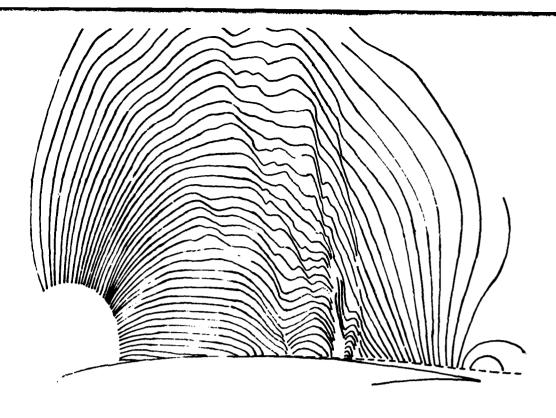
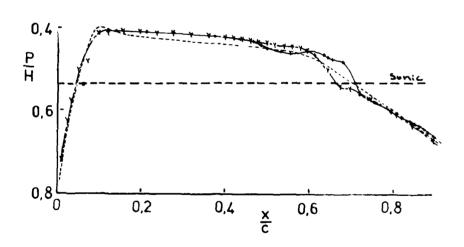


FIG. 11 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.780, $\alpha = 1.7^{\circ}$, TRANSITION FREE





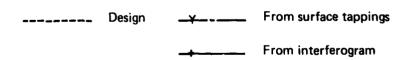
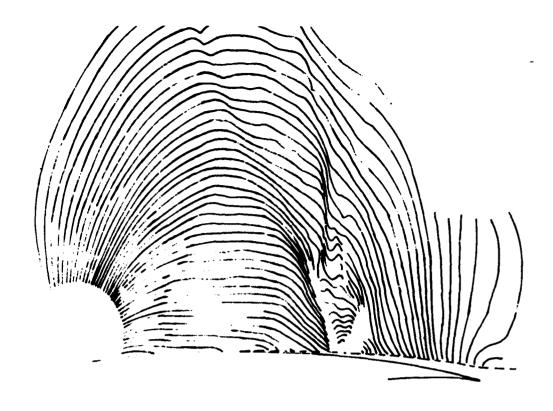


FIG. 12 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS $M=0.788,\,\alpha=1.7^{\circ},\,TRANSITION$ FREE



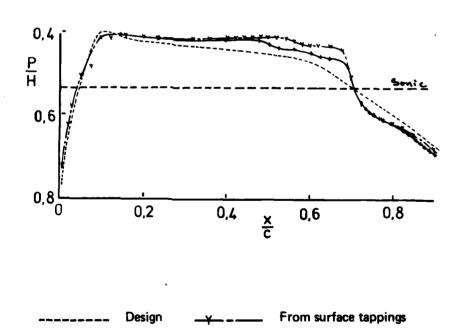
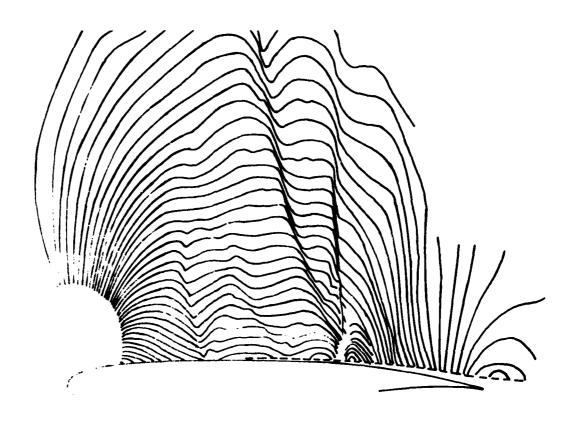


FIG. 13 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.793, α = 1.7°, TRANSITION FREE

From interferogram



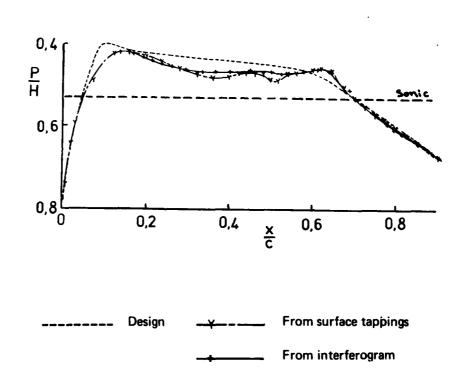
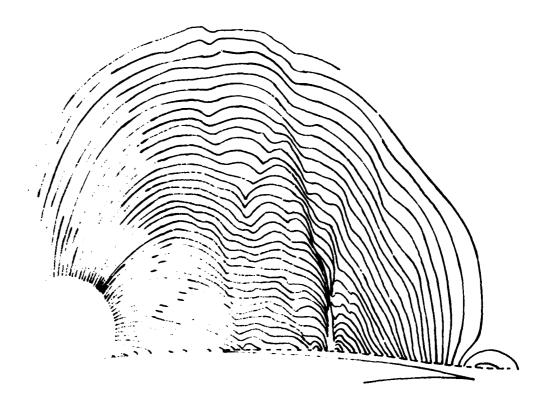


FIG. 14 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.2°, TRANSITION FREE



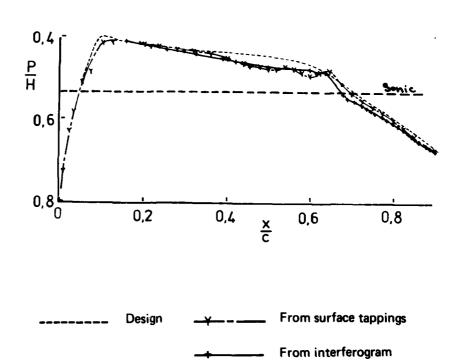
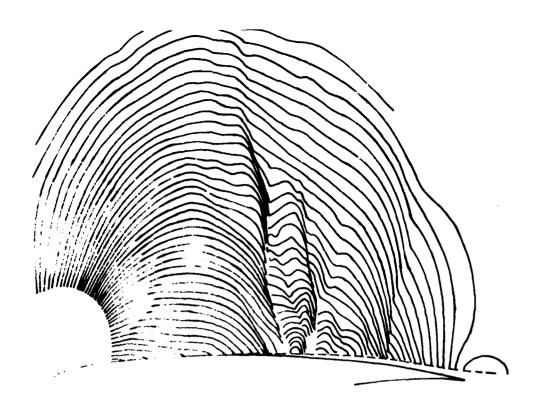
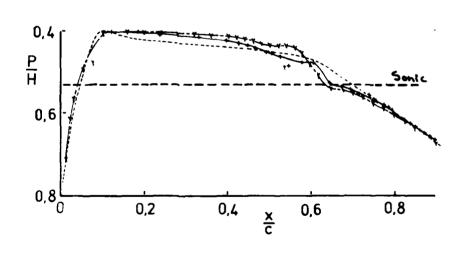


FIG. 15 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.5°, TRANSITION FREE





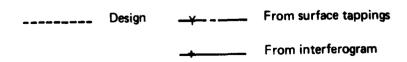
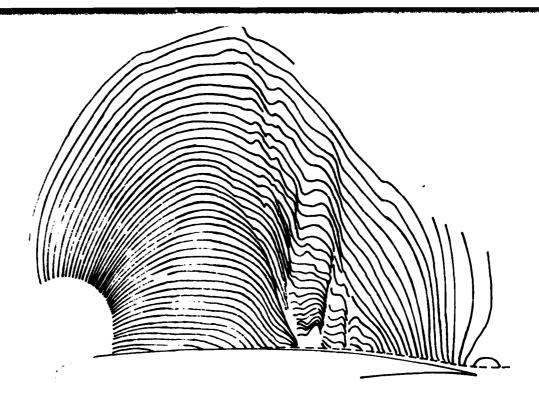
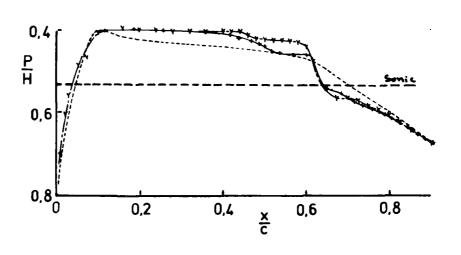


FIG. 16 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.9°, TRANSITION FREE





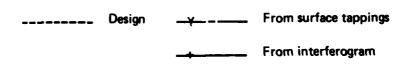
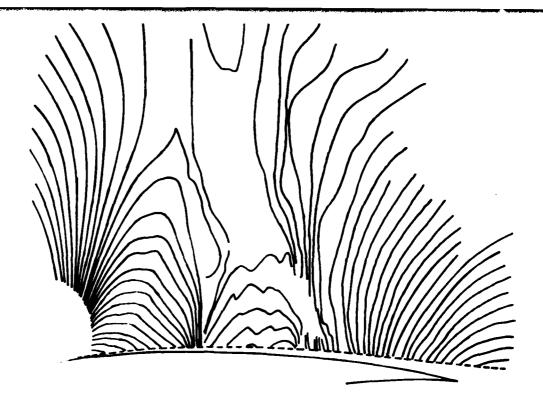
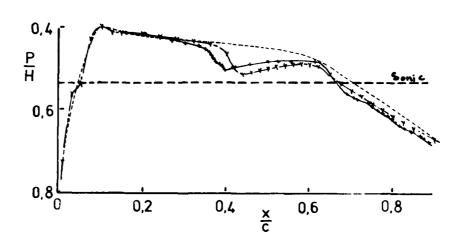


FIG. 17 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 2.1°, TRANSITION FREE





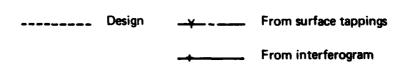
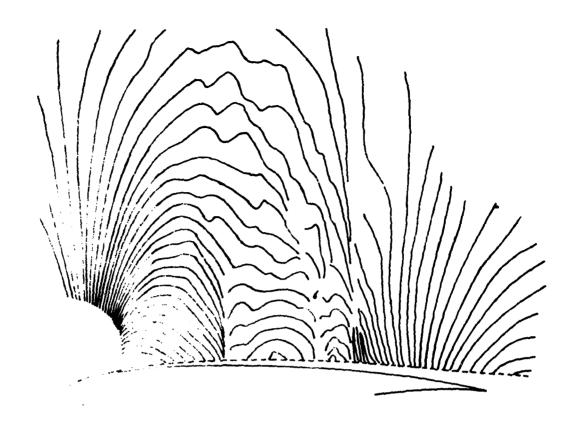
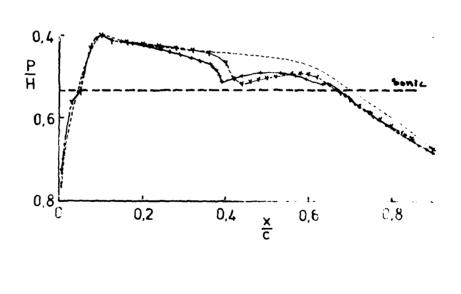


FIG. 18 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.7°, TRANSITION AT 2%C





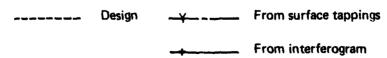
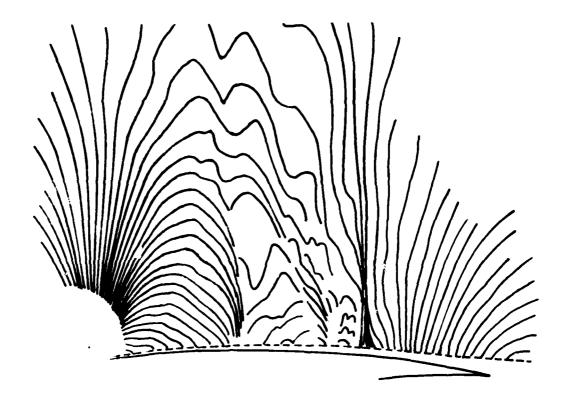
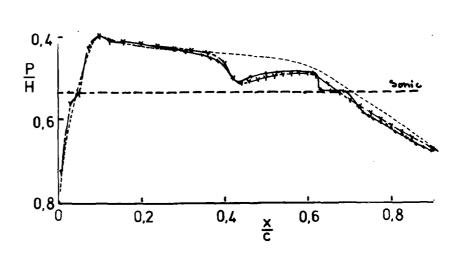


FIG. 19 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, $\alpha = 1.7^{\circ}$, TRANSITION AT 2%C





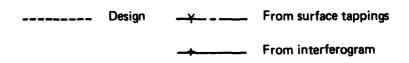
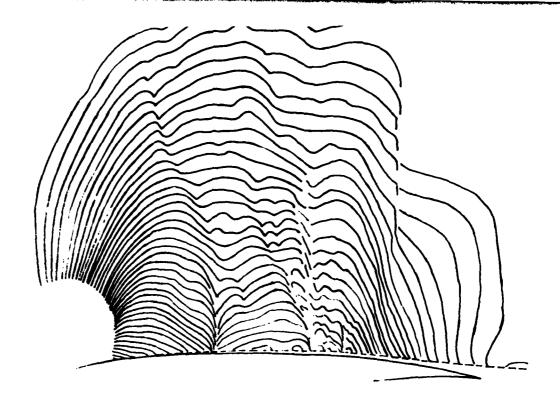
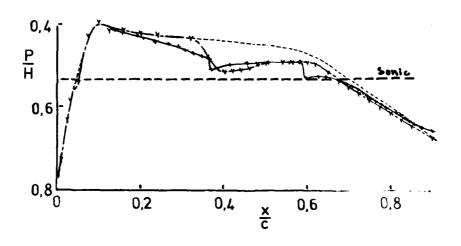


FIG. 20 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.7°, TRANSITION AT 2%C





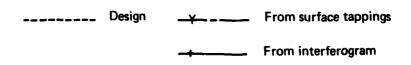
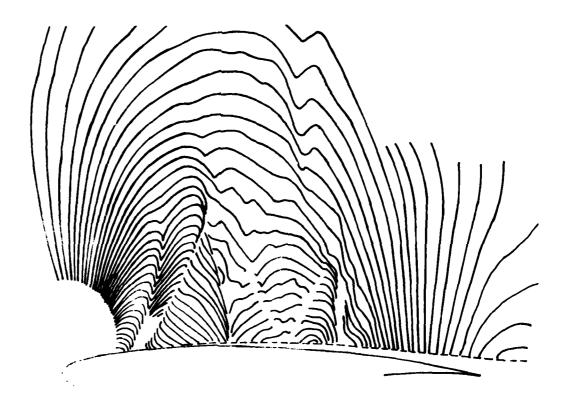


FIG. 21 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, $\alpha = 1.7^{\circ}$, TRANSITION AT 5%C





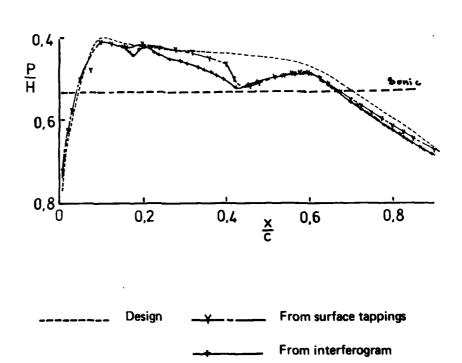
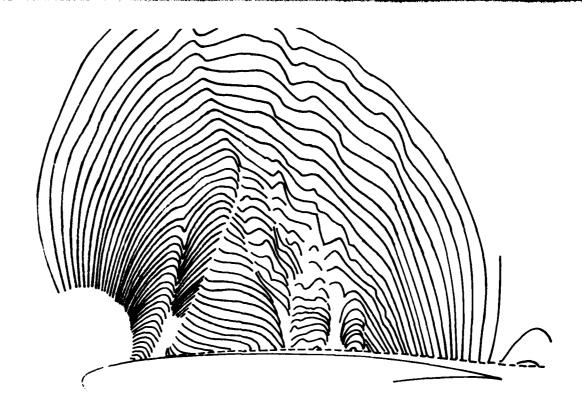
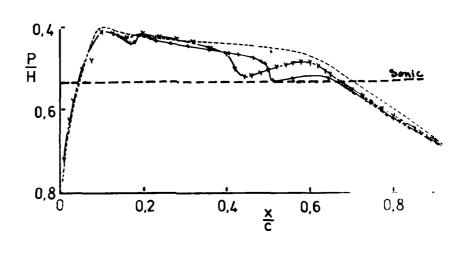


FIG. 22 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.7°, TRANSITION AT 17%C





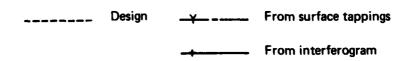
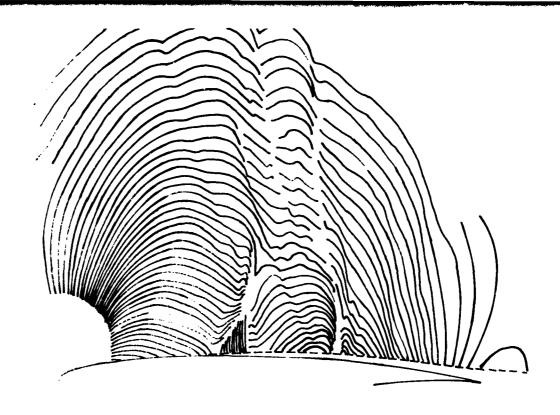


FIG. 23 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS $M=0.783, \alpha=1.7^{\circ}, TRANSITION AT 17\%C$



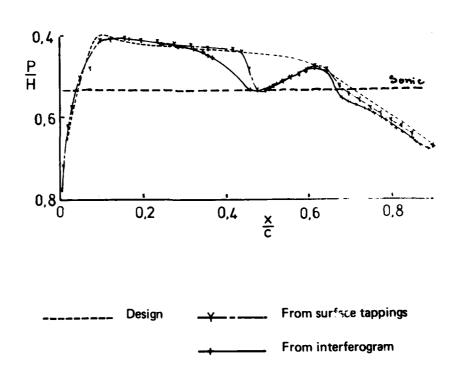
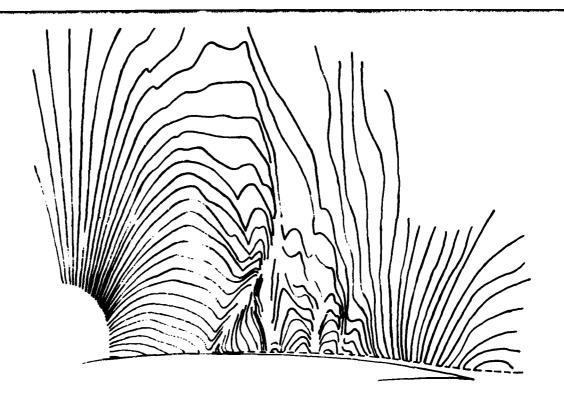
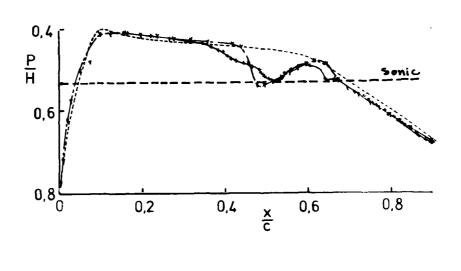


FIG. 24 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, α = 1.7°, TRANSITION AT 37%C





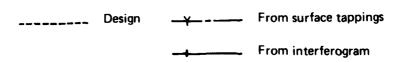


FIG. 25 INTERFEROGRAM AND PRESSURE DISTRIBUTIONS M = 0.783, $\alpha = 1.7^{\circ}$, TRANSITION AT 37%C

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17. Imprint Aeronautical R	esearch Laborate	ories, Melbourne			
18. Document Serie Aerodynamics		19. Cost Code 54 7710	20. Type of F	Report and Period	

